the shortest yet reported for a metal-silicon bond.²⁹ In the ruthenium complexes the two Ru-P bonds involving phosphorus atoms which are mutually *trans* to one another are normal, the mean values being 2.393 (6) Å (X = Cl) and 2.345 (4) Å (X = H), whereas the third Ru-P distance, involving the phosphorus atom which is *cis* to the other two, is in each case much shorter, being 2.230 (8) Å (X = Cl) and 2.206 (4) Å (X = H).^{8,9} These results may be rationalized if one assumes that transition metal-phosphine and --silyl bonds involve some back-donation from the metal; short metal-phosphorus and metal-silicon distances

(29) Further discussion of metal-silicon bond lengths appears in the following paper: Lj. Manojlović-Muir, K. W. Muir, and J. A. Ibers, *ibid.*, **9**, 447 (1970).

occur when there is no competing π -bonding ligand in the *trans* position. Since the covalent radius of phosphorus is some 0.07 Å shorter than the covalent radius of silicon, while the covalent radii of Ru and Rh differ by about 0.01 Å,³⁰ it would then appear, from the near equality of the short Ru–P distance in RuHCl(P-(C₆H₅)₈)₂ with the Rh–Si distance found here, that back-bonding occurs more readily with SiCl₈ than with P(C₆H₅)₃. Such an explanation is consistent with the relative stabilities of these complexes formed with different substituted silanes.¹²

Acknowledgment.—We are grateful to the National Science Foundation for its support of this study.

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The Crystal and Molecular Structure of Hydridobis(trichlorosilyl)carbonyl- π -cyclopentadienyliron, (π -C₅H₅)FeH(SiCl₃)₂(CO)

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The crystal structure of hydridobis(trichlorosilyl)carbonyl- π -cyclopentadienyliron, (π -C₃H₃)FeH(SiCl₃)₂(CO), has been determined by X-ray diffraction methods. A monoclinic unit cell of dimensions a = 7.493 (5) Å, b = 11.867 (9) Å, c = 8.651 (5) Å, and $\beta = 99.73$ (2)° contains two monomeric molecules; the observed and calculated densities are, respectively, 1.88 (2) and 1.835 g cm⁻², and the space group is C₂²-P2₁. The structure has been refined by least-squares methods to a conventional *R* factor of 0.041 over 1245 counter-diffractometric data. The absolute configuration of the molecules in the crystal used in the analysis has been determined by least-squares methods and Bijvoet ratio tests. The coordination polyhedron around the Fe atom can be regarded as a distorted tetragonal pyramid, with the centroid of the cyclopentadienyl ring at the apex and the remaining ligands in the basal plane. The Fe atom is displaced from the basal plane toward the ring. The two Fe–Si distances are equal at 2.252 (3) Å. This is the first determination of the length of such a bond. Structural data on transition metal-silicon bonds are summarized and arguments are presented in favor of some d_{π} - d_{π} back-bonding being involved in these bonds.

Introduction

Although complexes containing metal-silicon bonds are known for a wide variety of transition metals, there is still very little structural information on such compounds. The recent work on silylcobalt complexes of the type $Co(SiX_3)(CO)_4^{1-3}$ (where X = Cl, F, and H) tends to confirm the suggestion of Hagen and McDiarmid⁴ that the empty silicon d orbitals may be involved in the Co–Si bonding. We have pointed out in the previous paper that the Rh–Si distance in RhHCl(SiCl₃)-(P(C₆H₅)₈)₂ provides evidence that the silicon d orbitals are used in the Rh–Si bonding.⁵ In order to obtain further insight into the behavior of substituted silyl groups as ligands in transition metal complexes, we have determined the structure of π -CpFeH(SiCl₃)₂(CO) $(Cp = cyclopentadienyl = C_{\delta}H_{\delta}^{-})$. This is the first reported structure of a compound which contains an Fe–Si bond. The work also reflects our interest in the geometries of five-coordinated⁶ transition metal complexes and, in particular, those containing a hydridic hydrogen bonded to the transition metal.

Collection and Reduction of Intensity Data

A sample of π -CpFeH(SiCl₃)₂(CO) was kindly supplied by Graham.⁷ The crystals are pale yellow needles, moderately stable in air. A preliminary X-ray photographic examination showed that they possess monoclinic symmetry. According to the photographs of the 0kl, hk0, h0l, and h1l reciprocal lattice nets, taken with Mo K α radiation, systematic absences occur only

⁽¹⁾ W. T. Robinson and J. A. Ibers, Inorg. Chem., 6, 1208 (1967).

⁽²⁾ K. Emerson, P. R. Ireland, and W. T. Robinson, *ibid.*, 9, 436 (1970).

⁽³⁾ A. G. Robiette, G. M. Sheldrick, R. N. F. Simpson, B. J. Aylett, and

<sup>J. A. Campbell, J. Organometal. Chem. (Amsterdam), 14, 279 (1968).
(4) A. P. Hagen and A. G. McDiarmid, Inorg. Chem., 6, 686 (1967).</sup>

⁽⁵⁾ K. W. Muir and J. A. Ibers, ibid., 9, 440 (1970).

⁽⁶⁾ In this paper the cyclopentadienyl ring is regarded as occupying one vertex of the iron coordination polyhedron. Some authors assume that a cyclopentadienyl ring occupies three vertices of the coordination polyhedron and would therefore describe the iron as heptacoordinate in the present case.

⁽⁷⁾ W. A. G. Graham and W. Jetz, Abstracts, 155th National Meeting of the American Chemical Society, San Francisco, Calif., March-April 1968, No. M 082.

for the 0k0 reflections when k is odd. Such absences are compatible with the space groups C_{2h}^{1} -P2₁/m and C_2^2 -P2₁. The distribution of peaks in the threedimensional Patterson function indicated that the true space group is C_2^2 -P2₁; this assignment has been confirmed by a successful structure analysis. The unit cell dimensions, at 22°, are a = 7.493 (5) Å, b = 11.867(9) Å, c = 8.651 (5) Å, and $\beta = 99.73$ (2)°. These values were obtained by a least-squares procedure outlined below, and they are based on a value of 0.70930 Å for the wavelength of Mo $K\alpha_1$ radiation. The density of the crystals, measured by flotation in a mixture of carbon tetrachloride and dibromomethane, is 1.88 (2) g cm⁻³; it is in tolerable agreement with the value of 1.835 g cm⁻³ calculated for two molecules in the unit cell.

For the intensity measurements a large needle-shaped crystal was cut so that its final dimensions were 0.30 mm \times 0.54 mm \times 0.08 mm parallel to a, b, and c, respectively. It was then sealed in a glass capillary and mounted on a Picker four-circle diffractometer, in such a way that the b axis of the crystal made an angle of approximately 5° with the axis of the φ circle. The half-widths of a number of narrow source-open counter ω scans⁸ ranged from 0.08 to 0.21°, thus indicating that the crystal mosaicity was small, albeit somewhat anisotropic. Twelve reflections were accurately centered through a narrow vertical slit at a takeoff angle of roughly 0.5°. These observations were used to determine the crystal orientation and the dimensions of the unit cell by the method of least squares.⁹ The intensity measurements were made with the θ -2 θ scan technique, using procedures already described.¹⁰ Each reflection was scanned through a 2θ range of -0.7 to $+0.7^{\circ}$ from the center of the Mo K α_1 peak; the takeoff angle was 2° , and the scan rate was $1^{\circ}/\text{min}$. The background was counted for 10 sec at each end of the scan range. Molybdenum radiation was used in conjunction with a 3-mil niobium β filter and a pulse height analyzer set to accept roughly a 90% window centered on the Mo K α_1 peak. Coincidence losses were minimized by inserting copper foil attenuators into the diffracted beam until the maximum counting rate was less than 7000 counts/sec. The crystal-source and crystal-counter distances were 21 and 34 cm, respectively. The dimensions of the counter aperture were $5 \text{ mm} \times 5 \text{ mm}$.

The intensities of four strong independent reflections were measured periodically in order to monitor crystal and electronic stability. Although the crystal darkened during the experiment, no significant trend was apparent in the intensities of these standards. The greatest random fluctuation observed in the intensity of a standard was less than 3% of its mean value. The intensities were corrected for background, and their standard deviations were estimated in the way described earlier; a value for p of 0.04 was used.¹⁰ The data were then corrected for Lorentz and polarization effects. In order to correct for absorption, the six faces of the crystal were indexed on the X-ray diffractometer. All members of the forms $\{100\}, \{010\},$ and $\{001\}$ were exhibited. The dimensions of the crystal were measured with a micrometer eyepiece. The calculated transmission coefficients, based on a linear absorption coefficient of 21.9 cm⁻¹ for Mo K α radiation, ranged from 0.53 to 0.83.

The intensities of 1924 independent reflections within the sphere $2\theta(M \circ K \alpha_1) \leq 56^\circ$ were measured. Of these, 1251 were found to be greater than three times their estimated standard deviations. Only these reflections were used in subsequent Fourier and leastsquares calculations. Continuation past $2\theta = 56^{\circ}$ did not seem worthwhile, since in the shell with $2\theta >$ 52° there were only a few reflections with intensities significantly above background. For 1872 reflections only the *hkl* intensities were measured; for the remainder, which were all strong low-angle reflections, the intensities of hkl, \overline{hkl} , $h\overline{kl}$, $and \overline{hkl}$ were measured. After absorption corrections had been applied, the intensities of the *hkl* and *hkl* spectra were averaged, and so also were those of the *hkl* and *hkl* spectra, yielding the intensities of 52 Friedel pairs. These were later used to check the absolute configuration of the molecules in the crystal employed in the analysis. Reflections with negative k indices were not included in any of the least-squares calculations.

Structure Determination and Refinement

A three-dimensional Patterson synthesis revealed the positions of the iron, chlorine, and silicon atoms. The sites of the oxygen and carbon atoms were obtained from subsequent electron density syntheses. The model so derived was refined by the method of full-matrix least squares. The function minimized was $\Sigma w \Delta^2$, where $w = 1/\sigma^2(F_o) = 4F_o^2/\sigma^2(F_o^2)$ and where $\Delta = |F_o| - |F_c|$, $|F_o|$ and $|F_o|$ being the observed and calculated structure amplitudes, respectively. The atomic scattering factors were taken from ref 11, with the exception of those for the Fe¹² and H¹³ atoms. Throughout the refinement the anomalous scattering of the Fe, Cl, and Si atoms was taken into account in the calculation of structure factors.¹⁴ Cromer's values for $\Delta f'$ and $\Delta f''$ were used.¹⁵

A refinement in which the scale factor and the positional and isotropic temperature factors of all the nonhydrogen atoms were adjusted converged with R = 0.127 and $R_{\rm w} = 0.146$, where $R = \Sigma |\Delta| / \Sigma |F_{\rm o}|$ and $R_{\rm w} =$

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 - (15) D. T. Cromer, *ibid.*, **18**, 17 (1965).

⁽⁸⁾ T. C. Furnas, "The Single Crystal Orienter Instruction Manual," General Electric Co., Milwaukee, Wis., 1957.

⁽⁹⁾ Using the program PICK for the CDC 6400 computer. In addition to the various local programs, local versions of Hamilton's GONO absorption correction program, Busing and Levy's ORFFE function and error program, Johnson's ORFFE thermal ellipsoid plotting program, and Zalkin's FORDAP Fourier program were employed in this work. The local least-squares program NUCLS used in this work in its nongroup form closely resembles the Busing-Levy ORFLS program.

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⁽¹¹⁾ J. A. Ibers in "International Tables for X-Ray Crystallography," Vol. III, The Kynoch Press, Birmingham, England, 1962.

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(13) R. F. Stewart, E. R. Davidson, and W. T. Simpson, J. Chem. Phys.,

Observed and	CALCULATED	STRUCTURE	AMPLITUDES (IN	ELECTRONS X	\times 10) for π -CpF	eH(SiCl ₃) ₂ (CO)

TABLE I

 $(\Sigma w \Delta^2 / \Sigma w F_o^2)^{1/2}$. The poor agreement of this model is due mainly to the anisotropy of the atomic thermal vibrations. Adjustment of anisotropic thermal parameters for all the nonhydrogen atoms brought the agreement indices R and R_w to 0.046 and 0.053, respectively.

At this point the absolute configuration of the molecules in this particular crystal was considered. Let A be the configuration of the molecule assumed hitherto. The signs of the indices of all reflections were now reversed, thus assuming the enantiomorphic model B. Model B was refined in exactly the same way as model A. This refinement converged with R = 0.048 and $R_{\rm w} = 0.055$. Assuming that the data contain only random errors, B is a poorer model than A at the 0.5%significance level.¹⁶ The difference between $|F_{\circ}(h)|$ and $|F_{o}(\bar{h})|$ was greater than 2% of the mean $|F_{o}|$ for 15 of the 52 reflections for which Friedel pair intensities had been measured. Comparison of the ratio $|F_{\circ}(h)|/$ $|F_{o}(\bar{h})|$ with $|F_{c}(h)|/F_{c}(\bar{h})|^{17}$ again indicated that A was a better model than B; model A correctly predicted these ratios for all of the 15 reflections which showed Bijvoet differences of more than 2%, whereas model B

failed to predict any correctly. Furthermore, the agreement between chemically equivalent bond lengths was better for model A than it was for model B; thus the two Fe–Si distances differed by 0.006 Å (at this stage) for model A and by 0.027 Å for model B. These results are conclusive evidence that the correct absolute configuration is given by model A.

A difference map was now calculated, with the data for which $\lambda^{-1} \sin \theta < 0.30$ Å⁻¹, in order to locate the hydrogen atoms. The maximum electron density in this map was $0.32 \text{ e}^{-}/\text{Å}^3$, while the estimated standard deviation of the electron density $\sigma(\rho)^{18}$ was 0.08 e^{-/}Å³. Only one of the hydrogen atoms of the cyclopentadienyl ring was found to lie on a peak higher than $3\sigma(\rho)$; the remaining four lay in regions higher than $1.5\sigma(\rho)$, and the hydridic hydrogen atom could not be located from the map. Our failure to detect one hydrogen atom and the poor signal to noise ratio for the "peaks" associated with the five we did detect reflect the general experience that for Fourier syntheses in noncentrosymmetric space groups the peaks associated with atoms excluded from the structure factor calculations tend to be low in height. This failure does not seem to be due to a lack of sensitivity of our data to the hydrogen

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⁽¹⁷⁾ G. H. Stout and L. H. Jensen, "X-Ray Structure Determination," The Macmillan Co., New York, N. Y., 1968, p 426.

TABLE II FINAL FRACTIONAL COORDINATES AND ANISOTROPIC THERMAL PARAMETERS OF ATOMS

	T MAD T RECTAINED COORDINATES AND TRANSPIROPIC THERMAD TARAMETERS OF ATOMS								
	x	У	z	$\beta_{11}{}^a$	β_{22}	\$ 33	β_{12}	β_{13}	β_{28}
Fe	$0.1963(1)^{b}$	$\frac{1}{4^{\sigma}}$	0.2563(1)	1204(17)	627(9)	748(13)	-14(14)	106(12)	-1(13)
Si(1)	0.3726(3)	0.3921(3)	0.3663(3)	1619(46)	850(24)	986(32)	-224(27)	152(31)	-253(25)
Si(2)	0.0903(3)	0.2689(3)	-0.0019(2)	1901(43)	717(24)	734(26)	-65(27)	-21(27)	-37(23)
Cl(1)	0.2454(4)	0.5205(3)	0.4628(3)	3481(73)	926(26)	1716(43)	252(35)	241(44)	-499(28)
Cl(2)	0.5659(3)	0.3380(3)	0.5494(3)	2080(51)	1484(33)	1686(40)	-93(34)	-660(37)	-175(33)
Cl(3)	0.5236(4)	0.4715(3)	0.2232(3)	3565(78)	1621(40)	1629(42)	-1448(47)	806 (46)	-343(35)
Cl(4)	-0.1685(3)	0.3298(3)	-0.0570(3)	2299(55)	1258(32)	1787(43)	204(34)	-599(39)	313(32)
Cl(5)	0.2397(4)	0.3702(3)	-0.1231(3)	3731(74)	1241(31)	974(30)	-589(40)	545(38)	142(26)
Cl(6)	0.0808(4)	0.1173(3)	-0.1192(3)	4204(85)	946(29)	1399(40)	-65(38)	224(45)	-532(28)
0	0.5194(9)	0.1647(9)	0.1562(9)	2025(161)	1829(110)	2412(149)	536(109)	632(122)	-459(113)
С	0.3937(11)	0.2008(9)	0.1939(10)	1394(169)	1084(94)	1050(121)	13(101)	-60(117)	15(89)
C(1)	-0.0393(11)	0.1541(9)	0.2554(10)	1584(167)	962(94)	1249(128)	-294(102)	286(115)	107(98)
C(2)	-0.0426(11)	0.2493(11)	0.3495(11)	2014(164)	1103(93)	1729(142)	129(153)	1177 (130)	-87(133)
C(3)	0.1062(14)	0.2429(12)	0.4718(10)	3161 (240)	1322(120)	874(112)	-675(183)	584 (136)	-109(122)
C(4)	0,2028(14)	0.1451(11)	0.4506(12)	2674(240)	1052(113)	1371(162)	-14(134)	23(156)	559(117)
C(5)	0.1099(13)	0.0885(9)	0.3125(12)	2486(224)	850(95)	1601(165)	-80(118)	615(154)	-44(104)

^{*a*} The anisotropic temperature factor used was: $\exp[-(\beta_{11}h^2 + \beta_{22}k^2 + \beta_{33}l^2 + 2\beta_{12}hk + 2\beta_{13}hl + 2\beta_{23}kl)]$. The values of β_{ij} are multiplied by 10⁵. ^{*b*} Limits of error are standard deviations, given in units of the least significant digit of the quantity to which they refer. ^{*c*} This value is arbitrarily chosen to fix the origin.

scattering; when fixed contributions for the scattering of the hydrogen atoms of the cyclopentadienyl ring were added to the structure factors, further refinement resulted in agreement indices R = 0.043 and $R_w =$ 0.048. The positions of the hydrogen atoms were idealized for these calculations: the C-H distance was taken to be 1.0 Å: each hydrogen atom was assigned an isotropic thermal parameter B = 5.0 Å².

At this stage the six reflections for which $2\theta < 14^{\circ}$ were removed from the refinement; for these reflections $|F_{o}|$ was much greater than $|F_{c}|$, apparently because of errors in estimating the local background. After further calculations the refinement converged with R = 0.041 and $R_w = 0.044$. All parameters shifted by less than their estimated standard deviations in the last cycle of refinement. The standard deviation of an observation of unit weight was 1.26, close to the expected value of unity. The relative correctness of the weighting scheme was satisfactory, on the basis of the variation of the mean $w\Delta^2$ with $|F_o|$ and $\lambda^{-1} \sin \theta$. Extinction corrections did not appear to be necessary. The final difference map showed no important features; the largest function value was $0.47 \text{ e}^{-}/\text{Å}^{3}$, or 9% of the average height of a carbon atom in the final electron density synthesis. For none of the 673 reflections omitted from the refinements for which $F_0^2 < 3\sigma(F_0^2)$ did $|F_{o}^2 - F_{e}^2|$ exceed $3\sigma(F_{o}^2)$.

The final observed and calculated structure amplitudes for those reflections used in the refinements are listed in Table I. The final positional and thermal parameters of the nonhydrogen atoms are presented in Table II, and the calculated positional parameters of the hydrogen atoms are given in Table III.

Description of the Structure

The crystal structure of π -CpFeH(SiCl₃)₂(CO) is built by packing of discrete monomeric molecules. An inspection of all intermolecular contacts shorter than 3.9 Å reveals none which is appreciably smaller than the sum of the corresponding van der Waals radii. No intermolecular contacts are therefore listed here.

TABLE III CALCULATED FRACTIONAL COORDINATES OF HYDROGEN ATOMS

	x	У	z
$H(1)^a$	-0.1336	0.1364	0.1578
H(2)	-0.1352	0.3136	0.3335
H(3)	0.1431	0.3017	0.5567
H(4)	0.3167	0.1171	0.5189
H(5)	0.1457	0.0150	0.2724

 a Each hydrogen is marked with the same number as the carbon atom to which it is bonded.

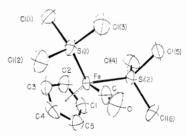


Figure 1.—A perspective view of the π -CpFeH(SiCl₃)₂(CO) molecule. The 50% probability ellipsoids of thermal motion are shown. Hydrogen atoms have been omitted for the sake of clarity.

A perspective view of the structure is shown in Figure 1 (from which the located hydrogen atoms have been omitted for the sake of clarity). The coordination polyhedron around the iron atom is of an unusual type. Although it may be described with equal validity in other terms, we choose to regard it as a distorted tetragonal pyramid, with the centroid of the cyclopentadienyl ring (CCp) at the apex; the carbonyl and trichlorosilyl ligands then lie in the basal plane, with the trichlorosily ligands trans to one another. The iron atom is placed above the basal plane, in the direction toward the cyclopentadienyl ring, so that the CCp-Fe-C angle is 126° (Table IV) and the CCp-Fe-Si angles are, respectively, 118 and 119°; as a result of the Si-Fe-Si angle is 115°. Inspection of molecular models indicates that a hydrogen atom can be accommodated trans to the carbonyl in the basal plane and at a reasonable Fe-H distance, without impossibly short non-

TABLE IV SELECTED INTERMOLECULAR DISTANCES AND ANGLES

SELECTE	D INTERMOLECU	JLAR DISTANCES AND	ANGLES
Distance	Value, Å	Angle	Value, deg
Fe-Si(1)	2.252(3)	CCp-Fe-Si(1)	119.4
Fe-Si(2)	2.252(3)	CCp-Fe-Si(2)	118.1
Fe–C	1.758(9)	CCp–Fe–C	125.8
Fe−CCp ^a	1.718	Si(1)–Fe– $Si(2)$	115.3(1)
Fe-C(1)	2.099(8)	Si(1)-Fe-C	85.1(3)
Fe-C(2)	2.084(7)	Si(2)-Fe-C	84.4(3)
Fe-C(3)	2.089(8)	Fe-C-O	177.0(10)
Fe-C(4)	2.086(10)	Fe-Si(1)-Cl(1)	116.7(1)
Fe-C(5)	2.106(11)	Fe-Si(1)-Cl(2)	112.2(2)
Fe-C _{ring} (av) ^b	2.093(4)	Fe-Si(1)-Cl(3)	115.8(1)
Si(1)-Cl(1)	2.049(4)	Fe-Si(2)-Cl(4)	115.3(1)
Si(1)-Cl(2)	2.060(4)	Fe-Si(2)-Cl(5)	115.8(1)
Si(1)-Cl(3)	2.043(4)	Fe-Si(2)-Cl(6)	112.2(1)
Si(2)-Cl(4)	2.049(4)	Fe-Si-Cl (av)	114.7(8)
Si(2)-Cl(5)	2.048(4)	Cl(1)-Si(1)-Cl(2)	103.5(2)
Si(2)-Cl(6)	2.061(4)	Cl(1)-Si(1)-Cl(3)	104.3(2)
Si-Cl (av)	2.052(3)	Cl(2)-Si(1)-Cl(3)	102.7(2)
C-O	1.132(10)	Cl(4)-Si(2)-Cl(5)	104.6(2)
C(1)-C(2)	1.395(14)	Cl(4)-Si(2)-Cl(6)	103.9(2)
C(2)-C(3)	1.404(13)	Cl(5)-Si(2)-Cl(6)	103.7(2)
C(3)-C(4)	1.396(17)	Cl-Si-Cl (av)	103,8(3)
C(4)-C(5)	1.443(15)	C(5)-C(1)-C(2)	109.9(9)
C(5)-C(1)	1.383(13)	C(1)-C(2)-C(3)	107.8(10)
C-C (av)	1.404(10)	C(2)-C(3)-3(4)	108.0(10)
		C(3)-C(4)-C(5)	107.9(9)
		C(4)-C(5)-C(1)	106.3(10)
		C-C-C(av)	108.0(6)

^a CCp is the centroid of the cyclopentadienyl ring. ^b Standard deviations of averaged quantities are estimated from the range of the individual measurements.

bonded intermolecular or intramolecular contacts. Our results, therefore, do not contradict the chemical evidence for the composition of this compound.

The geometry about the metal found here shows a strong resemblance to that found in the complexes π - $CpMoR(CO)_3$ (R = C_2H_{5} ,¹⁹ C_3F_7 ²⁰). In the molybdenum compounds the ring centroid occupies the apex of a tetragonal pyramid, with the alkyl and carbonyl ligands in the basal plane. The metal is again displaced from the basal plane toward the cyclopentadienyl ring so that the CCp-Mo-B angles (B = R, CO)are in the range 110-120°, and the B₁-Mo-B₂ angles $(B_1 trans to B_2)$ are smaller than 135° . No other structural studies on mononuclear iron complexes with similar geometries are known to us.

The thermal motions of the atoms (Figure 1 and Table V) show no unexpected features. The trichlorosilyl groups appear to be oscillating about the Fe-Si vectors. The directions of maximum vibration of the ring carbon atoms are all approximately normal to the corresponding ring carbon-CCp vectors and in the plane of the ring. Libration of the ring about the Fe-CCp vector would explain these results. The maximum root-mean-square amplitudes of all the nonhydrogen atoms are, however, quite small, and there is no evidence for disorder of the cyclopentadienyl ring.

Discussion

The cyclopentadienyl ring is planar within experi-

	1.	ABLE V	
ROOT-N	Mean-Square An	MPLITUDES OF VIE	bration (Å)
ALO	ng Principal Ax	es of Thermal I	Ellipsoids
Atom	Min	Intermed	Max
Fe	0.165(1)	0.185(1)	0.212(2)
Si(1)	0.172(4)	0.211(3)	0.261(4)
m • (-)			

TADLE V

Si(1)	0.172(4)	0.211(3)	0.261(4)
Si(2)	0.161(3)	0.225(4)	0.239(3)
Cl(1)	0.194(4)	0.284(3)	0.330(4)
C1(2)	0.179(3)	0.310(3)	0.329(4)
C1(3)	0.204(3)	0.235(3)	0.416(4)
Cl(4)	0.185(3)	0.302(4)	0.323(4)
C1(5)	0.176(3)	0.267(3)	0.351(4)
C1(6)	0.176(4)	0,296(4)	0.345(4)
0	0.202(10)	0.291(9)	0.384(11)
С	0.177(12)	0.221(12)	0.278(12)
C(1)	0.189(12)	0.216(11)	0.275(13)
C(2)	0.159(12)	0.283(12)	0.289(11)
C(3)	0.170(12)	0.247(15)	0.348(16)
C(4)	0.177(15)	0.271(13)	0.313 (15)
C(5)	0.228(12)	0.244(14)	0.272(12)
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mental error. The equation of the least-squares plane²¹ through the ring carbon atoms (based on monoclinic coordinates) is: 4.900X + 5.994Y - 5.758Z =-0.734. The deviations of the atoms from this plane are (Å): C(1), -0.005 (8); C(2), 0.007 (9); C(3), -0.006 (10); C(4), 0.003 (10); C(5), 0.003 (10). The ring C–C distances range from 1.383 (13) to 1.443 (15) Å, with a mean of 1.404 Å. On the basis of a χ^2 test the deviations of the individual values from the mean are not significant at the 1% probability level. Our results, therefore, do not suggest that localization of the ring electron density had occurred, although such localization has been reported in Cp-metal complexes where the electron density round the metal lacks effective cylindrical symmetry.²² The mean ring C-C distance is slightly shorter than that found in electron diffraction studies of ferrocene.23,24 This difference is probably due to our failure to correct the bond distances for librational motion of the ring rather than to electronic factors. The ring carbon-Fe distances range from 2.084 (7) to 2.106 (11) Å and are therefore equal within the experimental errors; their mean value is 2.093 (4) Å, and the Fe–CCp distance is 1.718 Å. The metal-cyclopentadienyl bonding is therefore somewhat weaker than in ferrocene, where values of 2.03 and 2.06 Å were obtained for the Fe-C distance in two separate electron diffraction analyses.^{23,24} Such a change is to be expected when one of the two rings of a sandwich molecule is replaced by strongly π -accepting ligands.²⁵ The geometry of the Fe-C-O group is normal.

The geometries of the two trichlorosilyl ligands are extremely similar. The same distortions from tetrahedral symmetry are found in the coordination about each silicon atom. The mean of the six Si-Cl distances is 2.052 Å, and the individual values all lie within

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TABLE VI	
Distances (\AA) and Angles (deg) in Silvl-Transition Metal Complexes $M-SiX_a$	3 ^a

MSi						
Compound	Obsd	Calcd ^b	Si-X	∠ M–Si–X	$\angle X$ -Si-X	Ref
π -CpFeH(SiCl ₃) ₂ (CO)	2.252(3)	2.51	2.052(3)	115	104	С
$Co(SiH_3)(CO)_4$	2.381(7)	2.51				d
$Co(SiCl_4)(CO)_4$	2.254(3)	2.51	2.035	113	105	e
$Co(SiF_3)(CO)_4$	2.226(5)	2.51	1.51	115	104	f
$Mn(Si(CH_3)_3)(CO)_5$	2.497(5)	2.63	1.91(2)			g
$RhHCl(SiCl_3)(P(C_6H_5)_3)_2$	2.203(4)	2.48	2.054(6)	109 - 121	103	h
$\mathit{trans-PtCl}(SiCH_3(C_6H_5)_2)(P(CH_3)_2C_6H_5)_2$	2.29	2.46				i

^a Averaged values are given whenever possible. ^b See text for covalent radii used. ^c Present work. ^d Reference 3. ^e Reference 1. ^f Reference 2. ^e R. S. Hamilton and E. R. Corey, Abstracts, 156th National Meeting of the American Chemical Society, Atlantic City, N. J., Sept 1968, No. INOR 025. ^h Reference 5. ^e P. Kapoor, Doctoral Dissertation, University of Sussex, 1968.

 ± 0.009 Å of this mean. The two Fe–Si distances are equal at 2.252 Å.

The available structural data on transition metalsilyl complexes are summarized in Table VI. Prediction of the length of a metal-silicon single bond from covalent radii is rendered difficult by the sensitivity of bond lengths involving transition metals to electronegativity effects, and to the coordination number and oxidation state of the metal; the special problems posed by Fe²⁶ and Co²⁷ in this respect have been recently pointed out. For Fe we propose a covalent radius of 1.34 Å; this value is based on the $Fe-C_{sp}$ distances of 2.123 (15) and 2.11 (2) Å found, respectively, in $Fe_2(C_2H_2)_3(CO)_6^{28}$ and π -CpFe(σ -Cp)- $(CO)_{2}^{29}$ and on the Fe-Fe distance of 2.679 (3) Å found in the former compound, where the metal atoms are not linked by a three-center bond. The upper limit of Fe-Fe distances in a number of iron-cluster compounds³⁰⁻³² is around 2.70 Å. The Co-Co distances in $[(CH_3CN)_5Co]_2^{2+33}$ and $Co_2(CO)_6(P(C_4H_9)_3)_2^{34}$ of 2.74 and 2.62 Å, respectively, lead to a Co radius of 1.34 Å. The Rh-CH₃ distance of 2.08 Å in RhI₂(CH₃)(P- $(C_6H_5)_3)_2^{35}$ gives an Rh radius of 1.31 Å. For Mn a value of 1.46 Å has been proposed from the Mn-Mn distance in Mn₂(CO)₁₀.³⁶ The Pauling radius for platinum is 1.29 Å.37 The metal-silicon single-bond distances predicted from these radii and from the accepted value of 1.17 Å for the covalent radius of silicon³⁷ are also given in Table VI.

From Table VI it can be seen that the observed metal-silicon distances are all shorter than their pre-

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dicted values, the differences ranging from 0.13 to 0.28 Å. Errors in the covalent radii used do not seem an adequate explanation for these differences, and we therefore conclude that some electronic effect is involved and that the shortening cannot be explained entirely by σ bonding. The shortening appears to be greatest when the silvl ligand carries electronegative substituents such as Cl and F; however, from the near equivalence of the Co-Si distances in Co(SiF₃)(CO)₄ and $Co(SiCl_3)(CO)_4$ it appears that the shortening is not purely a function of electronegativity. It would therefore seem that some $d_{\pi}-d_{\pi}$ back-bonding from the metal is involved and that the shortening is not entirely due to contraction of the metal d orbitals by a strongly electron-withdrawing ligand. Owing to a lack of data and to inadequacies in the values for the covalent radii, no firm conclusion is possible regarding trends in metal-silicon bonds as a given transition series is ascended. Yet consistent with the conjectured π -bonding effects, the shortest metal-silicon bond yet found involves the second transition series element rhodium in a structure where there is no π bonding ligand *trans* to and competing with the silvl group.

From Table VI it can be seen that the mean siliconhalogen distances in the complexes differ little from the values of 2.02 and 1.56 Å found for Si–Cl and Si–F distances in halogenosilanes.⁸⁸ The coordination polyhedron about the silicon atom in the metal complexes is a distorted tetrahedron; the metal–Si–X angle (X = Cl, F, etc.) is usually some 5° greater than the tetrahedral angle, whereas the X–Si–X angles are typically about 5° less than 109°. Similar distortions have been repeatedly noticed in the coordination about tin in transition metal–stannyl complexes.^{39,40}

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